

SYMMETRY - MOTIVATED

TUNING EFFECT IN PARTICLE MASSES

AND NUCLEAR DATA

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INTRODUCTION

- The Standard Model is a general theory of all interactions in which all three vector interactions (strong, weak and electromagnetic) are united in the representation:
- $$SU(3)_{\text{col}} \otimes SU(2)_L \otimes U(1)_Y$$
- Particle masses are parameters of the Standard Model. Y. Nambu suggested (1998) that empirical relations in particle masses could be used for the development of the Standard Model.
- The observed ratios are: $(m_\mu + m_e)/2(\delta m_\pi - m_e) = 13.00$. The lepton ratio $L=207=13 \times 16 - 1$. neutron mass $m_n+m_e/2(\delta m_\pi-m_e) = 115.007$. Ratio $m_n/m_e=1838.6836605(11)$, shift $\delta m_n = 161.65(6)$ keV.

$$\text{Ratio } \delta m_N / \delta m_n = 8 \times 1.0001(4).$$

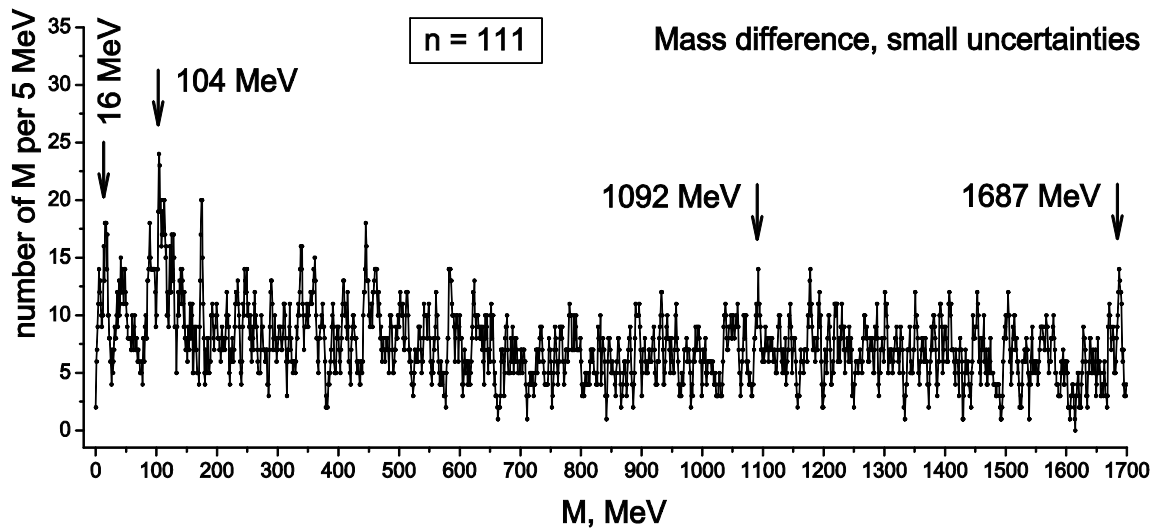
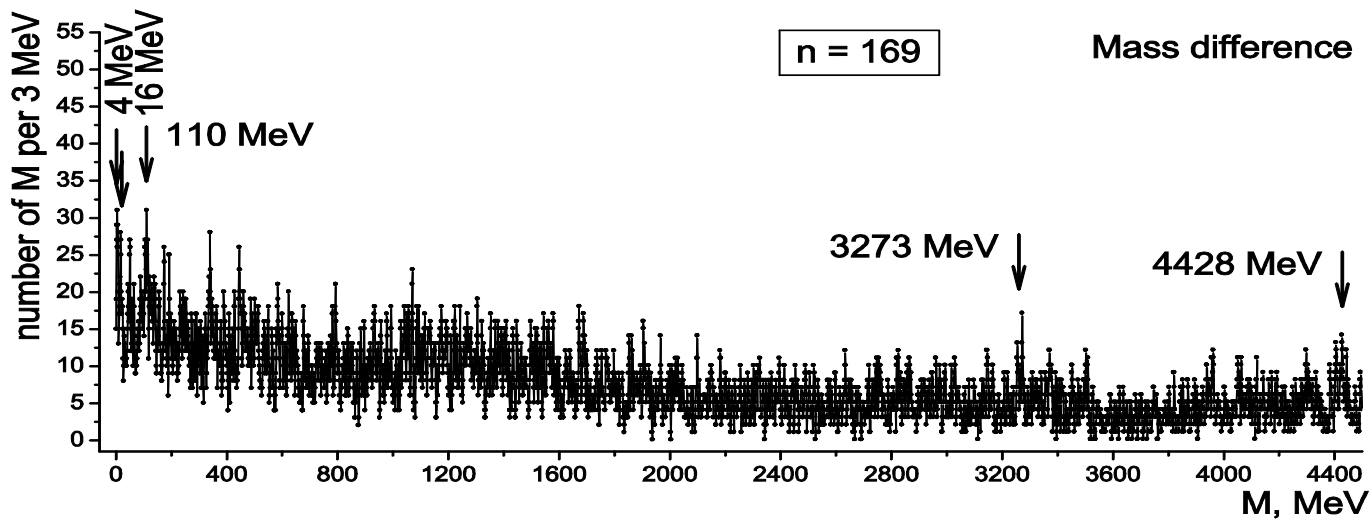
$$m_n - m_p = \delta m_N = 1.2933322(4) \text{ MeV};$$

$$m_n = 115 \times 16 m_e - m_e - \delta m_N / 8;$$

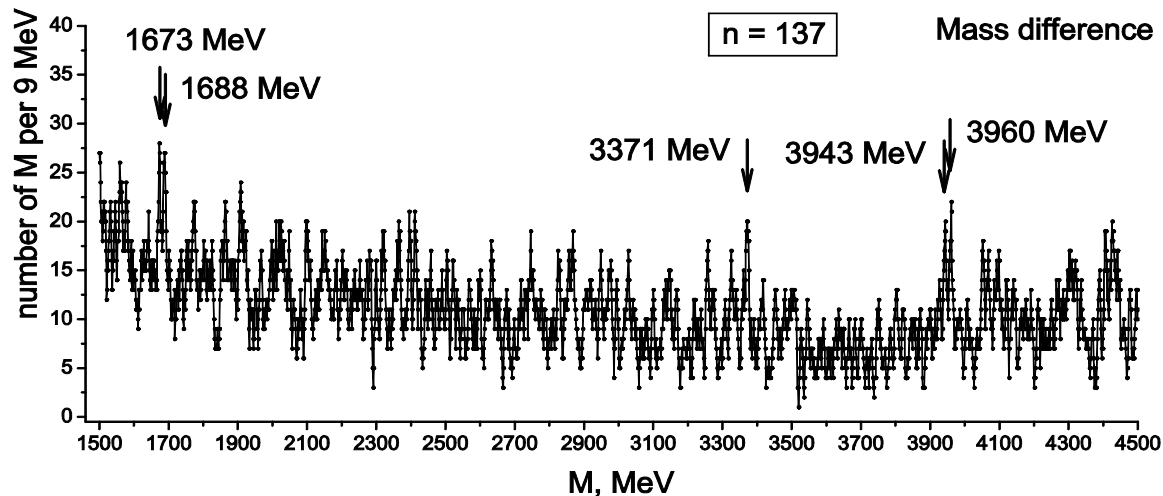
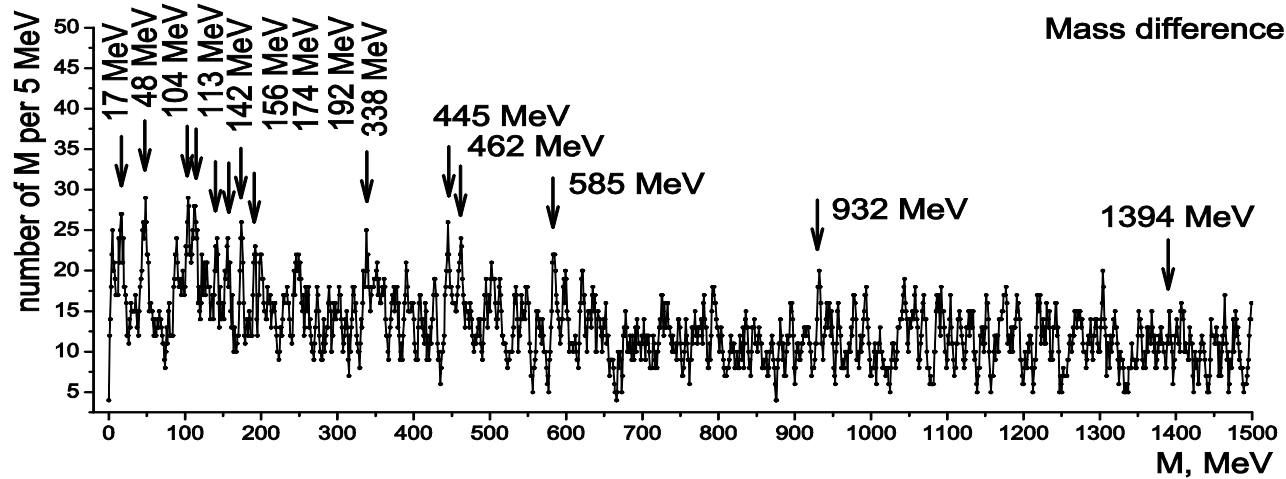
$$m_p = 115 \times 16 m_e - m_e - 9 \delta m_N / 8.$$

The lepton ratio $L=m_\mu/m_e=206.77$ becomes integer $207=9 \times 23=13 \times 16 - 1$ after a small QED radiative correction applied to m_e (it becomes $m_\mu/m_e(1-\alpha/2\pi)=207.01$). The factor $\alpha/2\pi = 115.9 \times 10^{-5}$, the QED radiative correction to the magnetic moment of the electron (Schwinger term) coincides with the relative value $117(11) \times 10^{-5}$ of the deviation of $\delta m_\pi = 4.5936(5)$ MeV from $9m_e = 4.5990$ MeV.

- Ratio of parameters of the Standard Model, masses of the muon and Z-boson: $m_\mu/m_Z = 115.9 \times 10^{-5}$.
- Belokurov and Shirkov suggested (1991) that electron mass also contains this factor $\alpha/2\pi$.
- A common approach is based on R. Feynman remarks: "The theories about the rest of physics are very similar to the theory of quantum electrodynamics: they all involve the interaction of spin $1/2$ objects with spin 1 objects. Why are .. the theories of physics so similar in their structure?"



Distribution of differences between all particle mass values from PDG-2016 (*top*) and values known relatively accurate (*bottom*).



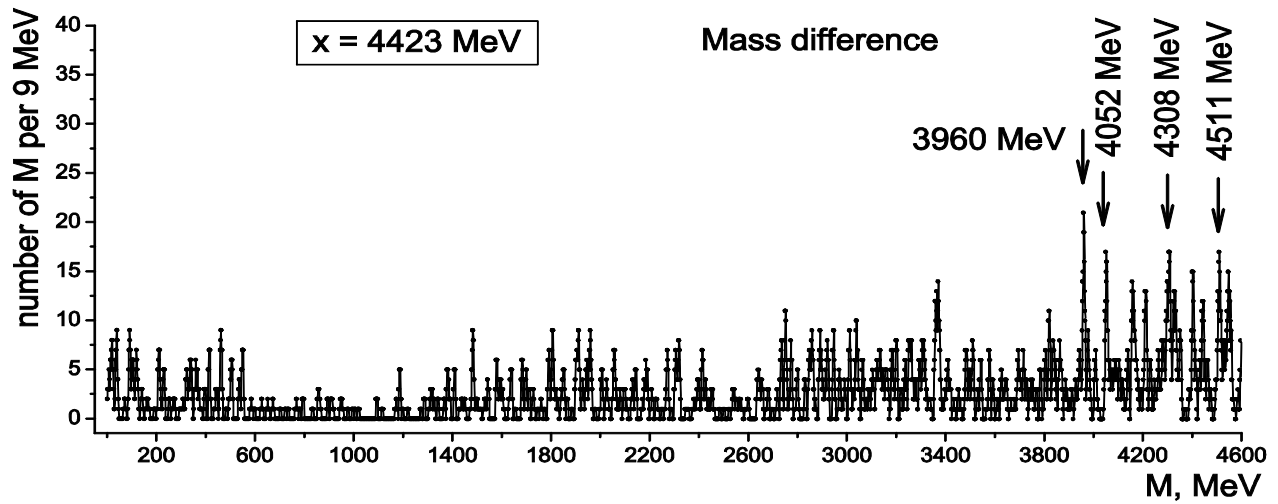
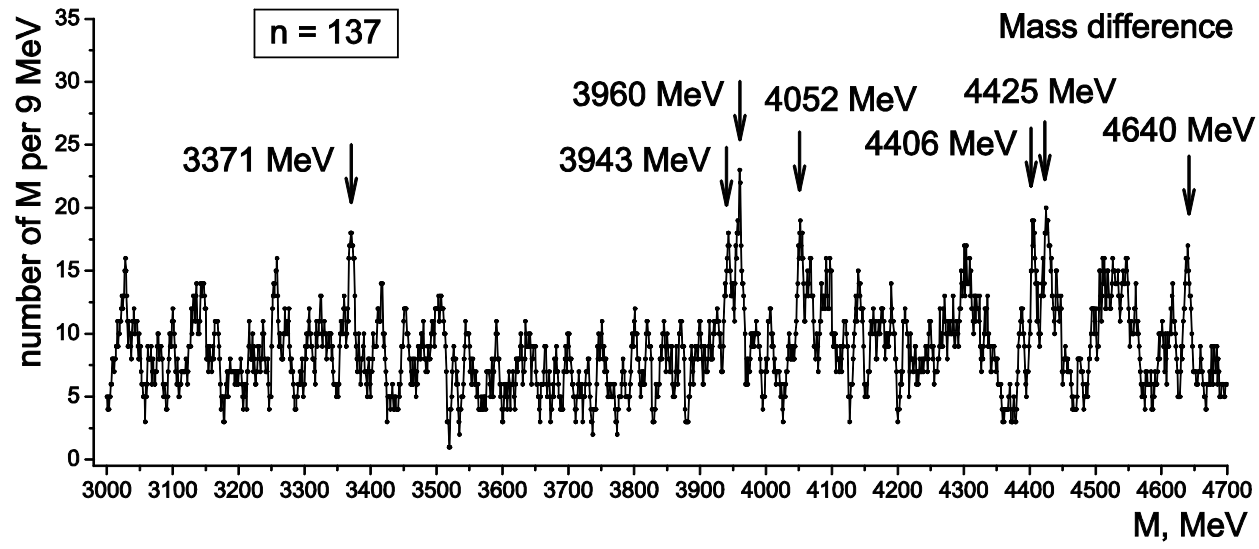
Differences between particle masses ΔM for averaging interval 5 MeV (*top*) and 9 MeV (*bottom*). Stable intervals (marked with arrows) are close to integer numbers of $\delta=16m_e$ found in CODATA relations (*top*) and to the b-quark mass estimation 4.2 GeV.

Table. Particle masses (in MeV) known with the uncertainty less than 6-10 MeV.

	Particle	m_1	Δ	17	48	104	142	156	174
1	leptons	electron, ν	0.0			106 (1)	140 (1)		
	μ		105.658			106 (1)			
	τ		1776.82		46 (15)	105 (7,8)			
2	Unflav.	mesons							
	f_π		130.7	0.4	***				
	π^0	$1^-(0^-)$	134.977		***				
	π^\pm	$1^-(0^-)$	139.570				140 (1)		
	η	$0^+(0^-)$	547.86						
	$\rho(770)$	$1^+(1^-)$	775.26						
	$\omega(782)$	$0^-(1^-)$	782.65					157 (1)	175 (1)
	$\eta'(958)$	$0^+(0^-)$	957.78		18 (1)			159 (2)	175 (1)
	$\phi(1020)$	$0^-(1^-)$	1019.46						173 (3)
	$b_1(1235)^*$	$1^+(1^+)$	1229.5	3.2		46 (2)		154 (4)	174 (4,5)
	$f_2(1270)^*$	$0^+(2^+)$	1275.5	0.8	19 (2)	46 (2,5)			
	$f_1(1285)$	$0^+(1^+)$	1282.0	0.5			102 (2)	142 (2,3)	
	$\eta(1295)^{**}$	$0^+(0^-)$	1294	4	19 (2)				
	$a_2(1320)$	$1^-(2^+)$	1318.3	0.5		46 (3)		138 (7)	
	$\eta(1405)^*$	$0^+(0^-)$	1408.8	1.8	18 (4,5)				
	$f_1(1420)^*$	$0^+(1^+)$	1426.4	0.9	18 (5)	47 (8,9)	105 (4,6)	142 (3)	154 (6)
	$\eta(1475)^{**}$	$0^+(0^-)$	1476	4		50 (7,9)		141 (5)	156 (7,8)
	$f_0(1500)^{**}$	$0^+(0^+)$	1504	6	16 (6)				158 (9)
	$f_2'(1525)^{**}$	$0^+(2^+)$	1525	5		49 (11)		142 (4,7)	
	$\pi_1(1600)^{**}$	$1^-(1^+)$	1662	8				142 (6)	158 (9)
$\eta_2(1645)^{**}$	$0^+(0^-)$	1617	5		50 (12)		141 (5)	156 (12)	
$\omega_3(1670)^{**}$	$0^-(3^-)$	1667	4		45 (12)		142 (7)	156 (13)	
$\pi_2(1670)^*$	$1^-(2^+)$	1672.2	3.0	17 (7)		105 (7)	140 (8)		
$\rho_3(1690)^*$	$1^+(3^-)$	1688.8	2.1	17 (7,8)				157 (11)	176 (6)
$f_0(1710)^{**}$	$0^+(0^+)$	1723	6		50 (13)		142 (10)		
$\phi_3(1850)^{**}$	$0^-(3^-)$	1854	7	16 (9)			142 (11)	156 (14)	
$a_4(2040)^{**}$	$1^-(4^+)$	1995	8	15 (10)	50 (17)		142 (11)		172 (7)
3	strange	mesons							
	K^\pm	$1/2(0^-)$	493.677						
	$K^*(892)^\pm$	$1/2(1^-)$	891.66			48 (1)			
	$K^*(892)^0$	$1/2(1^-)$	895.81		***				
	$K_1(1270)^{**}$	$1/2(1^+)$	1272	7		46 (3,4)		156 (3,5,6)	
	$K_1(1400)^{**}$	$1/2(1^+)$	1403	7	19 (3)				174 (4)
	$K_2^*(1430)^\pm$	$1/2(2^+)$	1425.6	1.5	17 (4)	47 (6,7)	104 (3,5)	142 (2)	154 (5)
	$K_2^*(1430)^0$	$1/2(2^+)$	1432.4	1.3	***				
	$K_2^*(1770)^{**}$	$1/2(2^-)$	1773	8				156 (12)	
	$K_3^*(1780)^{**}$	$1/2(3^-)$	1776	7		47 (14)			
	$K_4^*(2045)^{**}$	$1/2(4^+)$	2045	9		50 (17)			175 (8)
4	charmed	mesons							
	D^0	$1/2(0^-)$	1864.83			103 (9)	142 (10)		176 (6)
	D^\pm	$1/2(0^-)$	1869.58		16 (9)	47 (16)	141 (12)	155 (15)	175 (8)
	$D^*(2007)^0$	$1/2(1^-)$	2006.85		***				
	$D^*(2010)^\pm$	$1/2(1^-)$	2010.28		15 (10,11)		102 (10)	141 (12)	156 (14)
	$D_1(2420)^0$	$1/2(1^+)$	2420.8	0.5		50 (18,19)	103 (11)		157 (16)
	$D_2^*(2460)^0$	$1/2(2^+)$	2460.57		***				
	$D_2^*(2460)^\pm$	$1/2(2^+)$	2465.4	1.3			104 (12)		

Table.

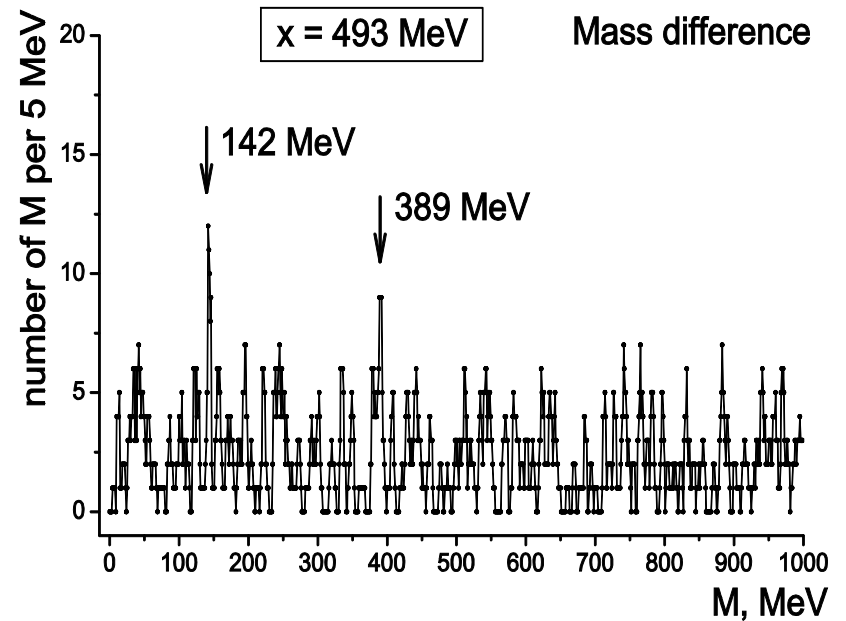
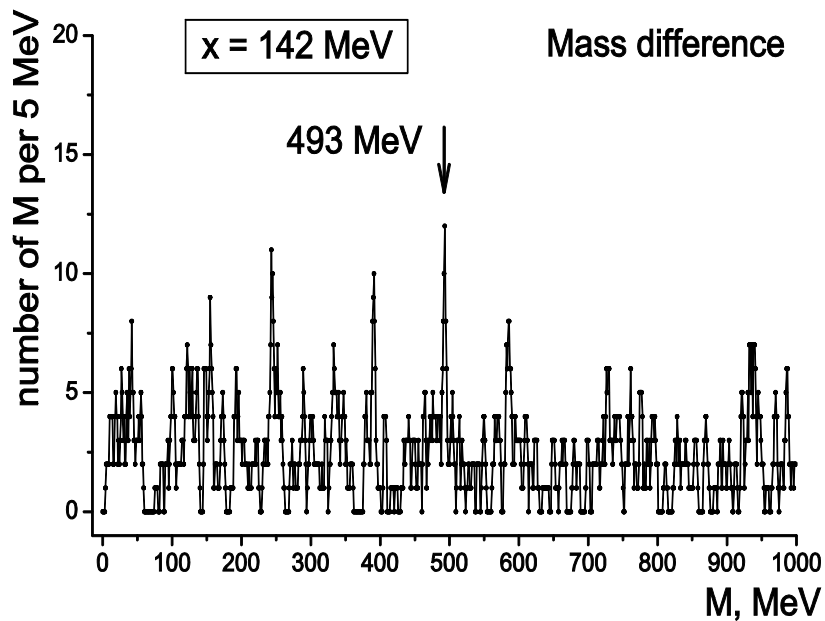
	Particle	m_1	Δ	17	48	104	142	156	174
5	charmed strange mesons								
	D_s^\pm	$0^+(0^-)$	1968.27			103 (9)	144 (13)		
	$D_s^*\pm$	$0(?)^?$	2112.1	0.4		102 (10)	144 (13,14)		174 (9)
	$D_{s0}^*(2317)^\pm$	$0(0^+)$	2317.7	0.6		103 (11)	142 (15)		
	$D_{s1}(2460)^\pm$	$0(1^+)$	2459.5	0.6			142 (15)		174 (10)
	$D_{s1}(2536)^\pm$	$0(1^+)$	2535.10		17 (14)				173 (13)
	$D_{s2}^*(2573)^*$	$0(2^+)$	2569.1	0.8			104 (12)		
$D_{s1}^*(2700)^\pm*$	$0(1^-)$	2708.3	3.4					173 (13,16)	
6	bottom mesons								
	B^\pm	$1/2(0^-)$	5279.31						
	B^0	$1/2(0^-)$	5279.62		***				
	B^*	$1/2(1^-)$	5324.65						
	$B_1(5721)^{+*}$	$1/2(1^+)$	5725.9	2.7			106 (17,18,19)		
	$B_1(5721)^{0*}$	$1/2(1^+)$	5726.0	1.3	***				
	$B_2^*(5747)^{+*}$	$1/2(2^+)$	5737.2	0.7			103 (20)		175 (23)
	$B_2^*(5747)^0$	$1/2(2^+)$	5739.5	0.7	***				
	$B_J(5970)^+$	$1/1(?)^?$ **	5964	5	15 (24)				172 (24)
	$B_J(5970)^0$	$1/1(?)^?$ **	5971	5	***				172 (24)
7	bottom strange mesons								
	B_s^0	$0(0^-)$	5366.82		49 (27)				
	B_s^*	$0(1^-)$	5415.4	1.5	49 (27)				
	$B_{s1}(5830)^0$	$0(1^+)$	5828.63		17 (22)		103 (18,21)		
$B_{s2}^*(5640)^0$	$0(2^+)$	5839.84			48 (28)	103 (20)			
8	bottom charmed mesons								
B_c^*	$0(0^-)$	6275.1	1.0						
9	$c\bar{c}$ mesons								
	$\eta_c(1S)$	$0^+(0^{+-})$	2983.4	0.5	15 (16)		102 (14)		
	$J/\psi(1S)$	$0^-(1^{--})$	3096.90		17 (17)			158 (18)	
	$\chi_{c0}(1P)$	$0^+(0^{++})$	3414.75				141 (18)		
	$\chi_{c1}(1P)$	$0^+(1^{++})$	3510.66		14 (18)	46 (22)			174 (21)
	$h_c(1P)$	$?^?(0^{+-})$	3525.38		14 (18)				
	$\chi_{c2}(1P)$	$0^+(2^{++})$	3556.20				46 (22)	141 (18)	
	$\eta_c(2S)^*$	$0^+(0^{+-})$	3639.2	1.2			47 (23)		
	$\psi(2S)$	$0^-(1^{--})$	3686.10				47 (23)		173 (21)
	$\psi(3770)$	$0^-(1^{--})$	3773.13				49 (24)	154 (19)	
	$\psi(3823)^*$	$?^?(2^{--})$	3822.2	1.2			49 (24)	105 (15)	
	$X(3872)$	$0^+(1^{++})$	3871.69		15 (19)	49 (25)			
	$X(3900)^*$	$1^+(1^{+-})$	3886.6	2.4	15 (19)				
	$X(3915)^*$	$0^+(?^{++})$	3918.4	1.9		47 (26)	106 (16)		
	$\chi_{c2}(1P)^*$	$0^+(2^{++})$	3927.2	2.6			105 (15)	154 (19)	
	$X(4020)^*$	$1(?)^?$	4024.1	1.9	15 (20)		106 (16)		
	$\psi(4040)^{**}$	$0^-(1^{--})$	4039	1	15 (20)				
	$X(4140)^*$	$0^+(?^{?+})$	4146.9	3.1					
	$\psi(4160)^{**}$	$0^-(1^{--})$	4191	5				156 (20)	
	$X(4260)^{**}$	$?^?(1^{--})$	4251	9					
$X(4360)^{**}$	$?^?(1^{--})$	4346.9	6				156 (20)		
$\psi(4415)^{**}$	$0^-(1^{--})$	4421	4						
$X(4660)^{**}$	$?^?(1^{--})$	4643	9						
10	$b\bar{b}$ mesons								
	$\eta_b(1S)^*$	$0^+(0^{+-})$	9399.0	2.3					
	$\Upsilon(1S)$	$0^-(1^{--})$	9460.30						
	$\chi_{b0}(1P)$	$0^+(0^{++})$	9859.44						
	$\chi_{b1}(1P)$	$0^+(0^{++})$	9892.78		19 (25)				
	$h_b(1P)^*$	$?^?(1^{+-})$	9899.3	0.8					
	$\chi_{b2}(1P)$	$0^+(2^{++})$	9912.21		19 (25)				
	$\Upsilon(2S)$	$0^-(1^{--})$	10023.26					140 (21)	
	$T(1D)^*$	$0^-(2^{--})$	10163.7	1.4			105 (23)	140 (21)	
	$\chi_{b0}(2P)$	$0^+(0^{++})$	10232.5	0.4					
	$\chi_{b1}(2P)$	$0^+(1^{++})$	10255.46						
	$\chi_{b2}(2P)$	$0^-(2^{+-})$	10268.65				105 (23)		
	$\Upsilon(3S)$	$0^-(1^{--})$	10355.2	0.5				157 (22)	174 (25)
	$\chi_{b1}(3P)^*$	$0^+(1^{++})$	10512.1	2.3	17 (26)			157 (22)	
	$\Upsilon(4S)^*$	$0^-(1^{--})$	10529.4	1.2	17 (26)				174 (25)
	$X(10610)^\pm*$	$1^+(1^+)$	10607.2	2.0					
	$X(10610)^{0**}$	$1^+(1^+)$	10609	6					



Top: Spacing distribution in mass spectrum in high-energy region.

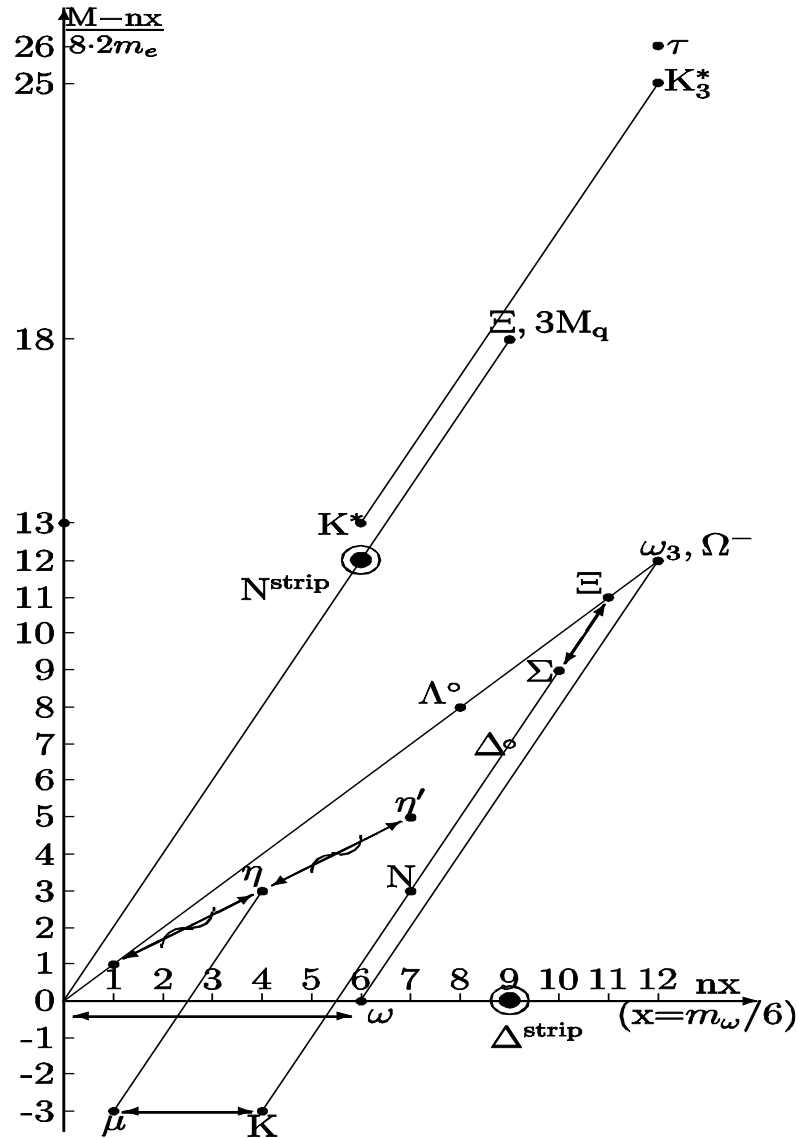
Bottom: Application of AIM Method to $x=4423$ MeV (AIM upward and downward directions, separately).

Distribution of intervals adjacent to $x=142$ MeV (m_{π}) and $x=493$ MeV (m_k).



Tuning effect in particle masses.

Different estimates of constituent quark masses.



Estimation of the baryon constituent quark mass $M_q=441$ MeV

Fig. 2a. (C.D. Roberts et al.). QCD gluon-quark-dressing effect calculated with Dyson-Schwinger Equation, initial masses m ; constituent quark mass arises from a cloud of low-momentum gluons attaching themselves to the current-quark; this is chiral symmetry breaking: nonperturbative effect that generates a quark mass from nothing even at $m=0$ (bottom).

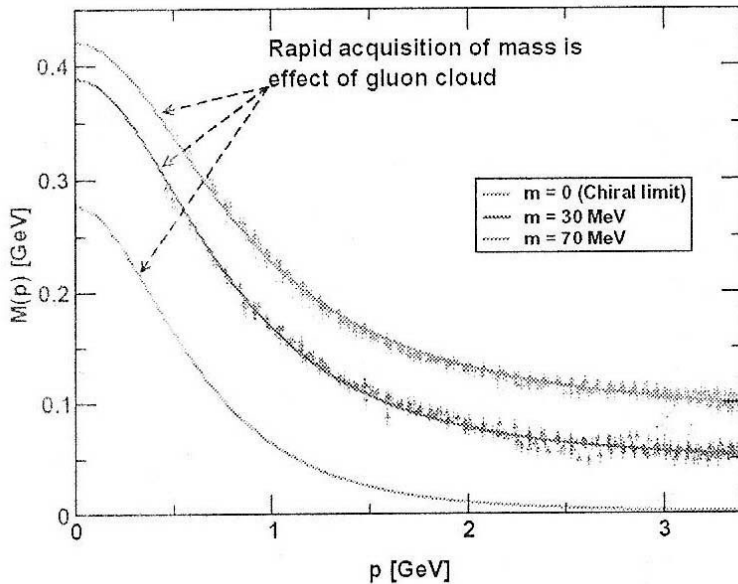


Fig.2b. (L.Glozman *et al.*). Calculation of nonstrange baryon (left) and lambda-baryon masses as a function of interaction strength within Goldstone Boson Exchange interaction Constituent Quark Model; initial baryon mass 1350 MeV = $3 \cdot 450$ MeV = $3M_q$ is marked as bottom “+” on left vertical axis.

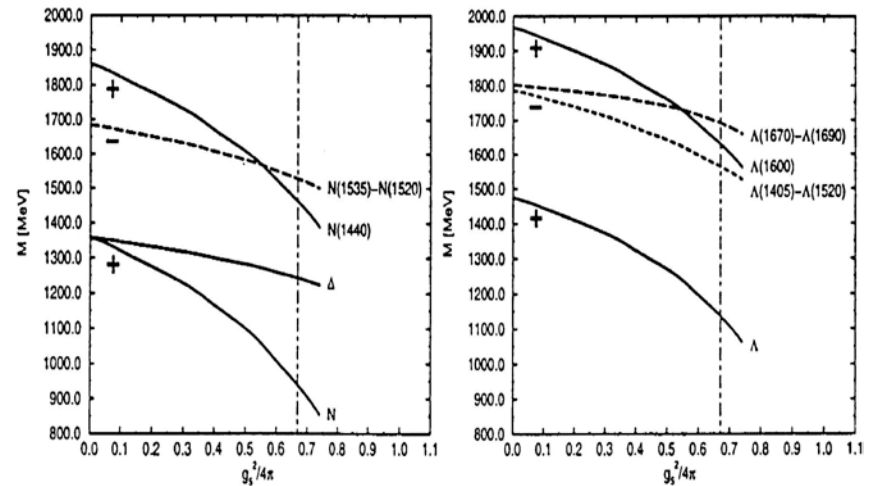
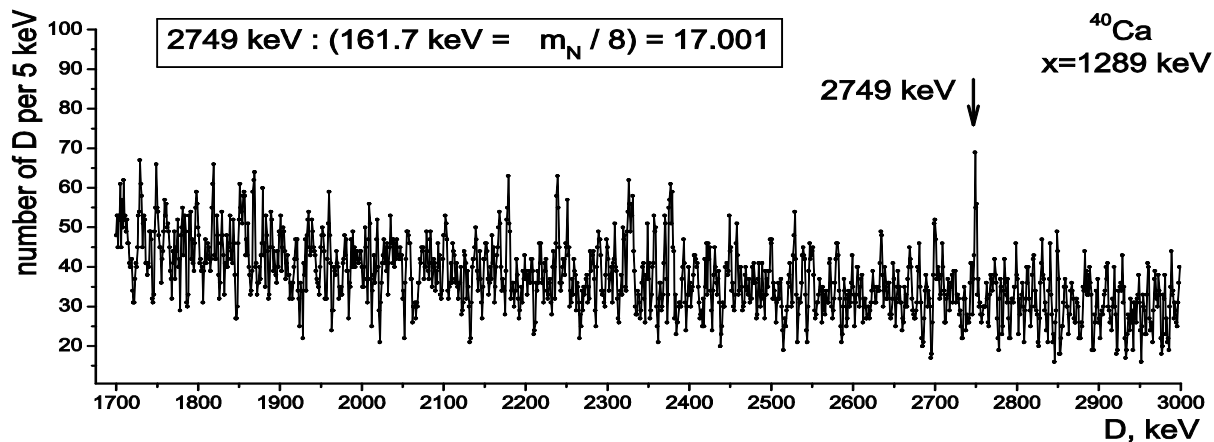
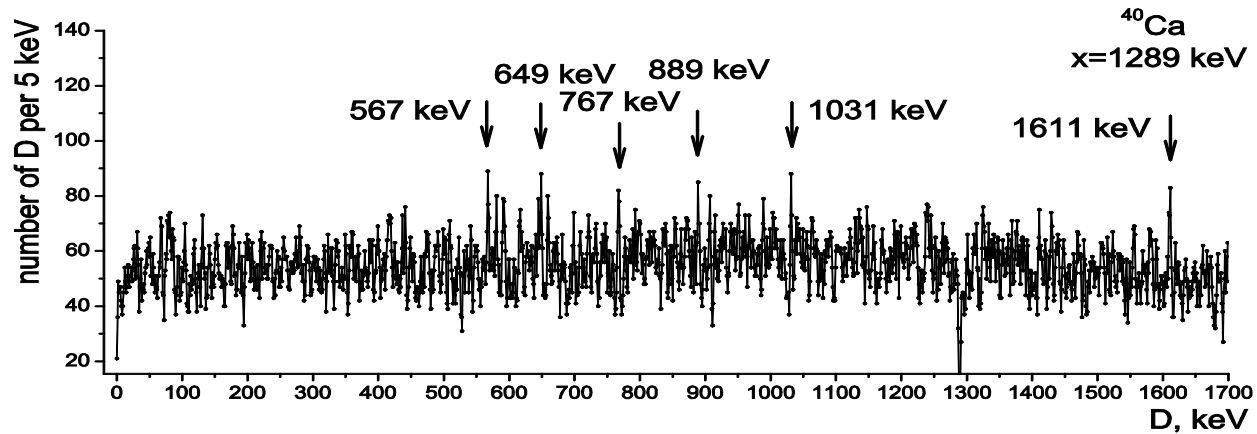


Table B. Linear dependence of excitations in near-magic nuclei (double boxed) upon numbers of protons in Z=14-20 shell (1st line) and numbers of valence neutrons (2nd line, boxed) is compared with the integer numbers of the parameter of 161 keV = $\Delta^{\text{TF}} = \delta m_N / 8$ determined in Z=50,51 nuclei (see Proc. QCD14, P.270, Fig.2 and Proc. QCD15, P.214, Table 3, Figs. 3-5).

$(Z-14)/2$	3	2	1	1	1	1	0	0	
N			$\Delta N=1$	$\Delta N=2$		$\Delta N=7$			
$A Z$	^{41}Ca	^{39}Ar	^{37}S	^{38}S	^{33}S	^{43}S	^{32}Si	^{35}Si	
E^*	0.0	1943	1267	646.2	1292	322	320.7	1942	973.9
$2J^\pi$	7^-	3^-	3^-	3^-	2^+	D	7^-	2^+	(3^+)
$n \frac{\delta m_N}{8}$	0.0	1941	1293	646	1293	322	322	1941	971
n		12		4	8	2	2	8	6
$A Z$	^{33}Mg	^{41}K		^{47}Sc	^{47}V	^{50}V	^{51}V	^{55}V	
E^*	159	484	980.4	1293.6	807.9	1294.9	320.2	320.1	323.3
$2J_o^\pi$	3^-		3^+		2^+	3^-	6^+	7^-	(7^-)
$2J^\pi$	(7^-)	(3^-)	4^+	7^-	3^-	11^-	4^+	5^-	(5^-)
$n \frac{\delta m_N}{8}$	161	483	971	1293	808	1293	322	322	322
n	1	3	6	8	5	8	2	2	2



D^{AIM} - distribution in ^{40}Ca levels for $x=1289$ keV close to $8 \times 17\delta' = \delta m_N = 1293.3$ keV in two energy regions (in upwards direction, $\Delta = 3$ keV) with the maximum at 2749 keV $\approx (17/8) \times 1293.3$ keV = 2748 keV.

Table C. New data (*) on ^{53}Ni and ^{53}Co excitations with the configuration of nucleon as three holes in ^{56}Ni core. Parameters of the residual interaction of valence nucleons and holes (double boxed) are compared with the same parameters in Z-50,51- and Z,N - nuclei. Boxed are excitations corresponding to maxima in sum distributions shown in Figs.6-8 in Proc. QCD-15.

AZ	$^{53}\text{Ni } 2J_o=7-$	^{58}Ni	^{59}Ni	^{61}Ni	^{63}Ni				
E^*	320(3)	1292*	1456*	1454.2	339.4	1454.8	87.1	1289.1	1451
$2J^\pi$	(5^-)	(3^-)	(11^-)	2^+	3^-5^-	7^-	1^-5^-	9^+	$(5,7,9)$
$n \frac{\delta m_N}{8}$	322	1293	1454	1454	322	1454		1293	1454
n	2	8	9	9	2	9		8	9
AZ	^{53}Mn				^{55}Mn				
E^*	378	1289.9	1441.3	2563.1	2573.1	1289.1	1292.1	1293.0	2582
$2J^\pi$	5^-	3^-	(11^-)	13^-	7^-	5^-11^+	11^-	(1^-)	
$n \frac{\delta m_N}{8}$	322	1293	1454	2586	2586	1293	1283	1293	2586
n	2	8	9	16	16	8	8	8	16
AZ	^{53}Co	^{59}Co				^{69}Cu	^{71}Cu		^{73}Cu
E^*	646.2*	1291.6	1459	2581.7	2585.8	1297.9	1453.3	2576(3)	1298.0
$2J^\pi$	7^-	3^-	11^-	3^-7^-	7^-	$3^-1, 3^-$	3^-9^-	(13^-)	$(3^-, 7^-)$
$n \frac{\delta m_N}{8}$	647	1293	1454	2586	2586	1293	1454	2586	1293
n	4	8	9	16	16	8	9	16	8

Table. Representation of parameters of tuning effects in particle masses (top) and nuclear data (bottom) with the expression $n \cdot 16m_e (\alpha/2\pi)^X M$ and different values of the X-power of QED factor $\alpha/2\pi$ and integers M and $n=1,13-18$. Boxed are five groups of values differing with $\alpha/2\pi=115.9 \cdot 10^{-5}$.

X	M	n = 1	n = 13	n = 16	n = 17	n = 18	n = 18.6	Comments
-1	3/2			$m_t=172.0$				
GeV	1	$16M_q=\delta^\circ$	$M_Z=91.2$	$M_H=115$		$M_H=126$		$\delta^\circ=7.06$
	1/2	(m_b-M_q)		$M^{L3}=58$				
0	1	$2m_d-2m_e$	$m_\mu=106$	$f_\pi=130.7$	$m_\pi-m_e$	$\Delta M_\Delta=147$	$2M_q$	
MeV	3			$M''_q=m_\rho/2$		$M_q=441=\Delta E_B$		NRCQM
1	1	$16m_e=\delta=8\varepsilon_o$	118		$k\delta-m_n-m_e=$	$170 = m_e/3$		Part.
keV	8				$=161.651$			mass
					$\delta m_N=1293$			CODATA
1	1	$9.5=\delta'=8\varepsilon'$	123	152	$\Delta^{TF}=161$	170 (Sn)	$\varepsilon_o=2m_e$	
keV	3				484 (E^*)	512 (Pd)		
	4		492		648 (Pd)	682 (Co)		Nuclear
	8		984	1212	1293 (E^*)	1360 (Te)		data
2	1,4	$11=\delta''=8\varepsilon''$	143	176	749 (Br,Sb)		$\varepsilon'=1188$	Neutron
eV	4,8		570 (Sb)		1500 (Pd,Hf)	X=3	$\varepsilon''=1.35$	reson.

Conclusions

Here we used data on scalar mass $M=125.0$ GeV and CODATA ratio m_n/m_e to find out an explanation of the systematic appearance of nuclear mass/energy intervals close to mass differences of the nucleon, leptons and the pion. The relation $(n \cdot 16m_e - m_e - m_n)/\delta m_N = 1/8.001(2)$ and the similar long-range correlations in particle masses could be checked with the new data.

- Relation between observed stable nuclear intervals and particle masses can be considered in connection with recently obtained mass of the scalar field, the ratios $m_\mu/M_Z = \alpha/2\pi$ and $(1/3m_e)/M_H = (\alpha/2\pi)^2$. There is a possibility that they are reflections of the fundamental relations between SM-parameters (which were mentioned as a future “super-duper” model by R. Feynman).
- Observed analogy between tuning effects in particle masses and in nuclear data should be theoretically based on QCD as a part of Standard Model. It is in line with Y. Nambu suggestion about the role of empirical relations in particle masses for the SM development.
- Scientific potential of nuclear physics can be connected with a fundamental role of QCD in the Standard Model and with the role of QED parameters.

Thank you for your attention